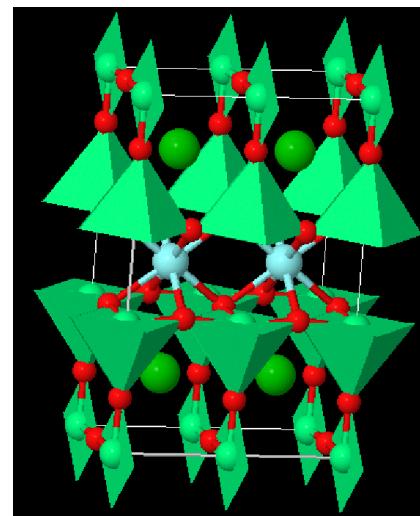




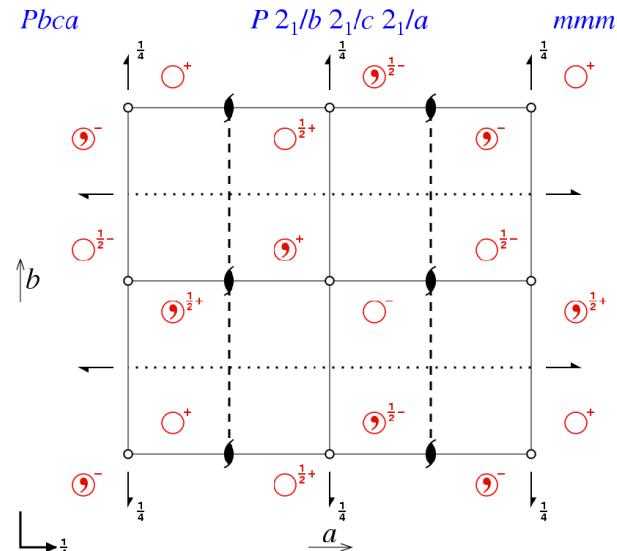
# Ecole du GDR MICO

(Matériaux et Interactions en COmpétition)

Mai 2014



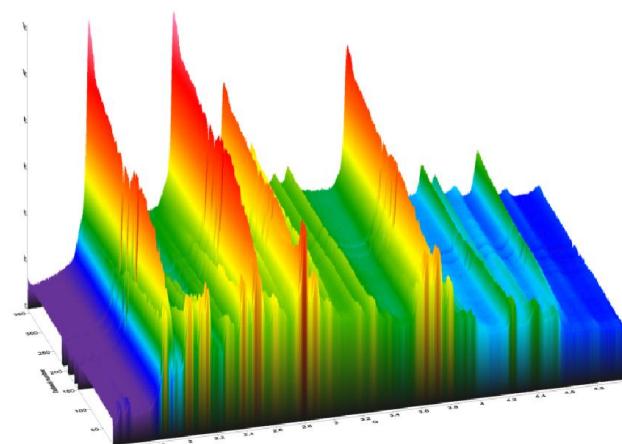
# Cristallographie et techniques expérimentales associées



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# Outline

## part I: CRYSTALLOGRAPHY IN DIRECT SPACE

### I.1. Orientation symmetry

Elementary point symmetries

How to obtain and name all crystallographic point groups?

Examples of point groups

The 32 point groups and 11 Laue classes

### I.2. Translation symmetry

Lattice and motif, Unit cell

The orientation symmetries of lattices:

the 6 conventional cells, 7 crystal systems and 14 Bravais lattices

Rows and net planes

### I.3. Space group symmetry

Glide planes and screw axes

The 230 space groups

The International Tables for Crystallography

### I.4. Beyond basic crystallography

Aperiodic crystals: superspace groups

Magnetic structures: magnetic point groups and space groups



# Outline

## part II: DIFFRACTION - CRYSTALLOGRAPHY IN RECIPROCAL SPACE

### II.1. The reciprocal space

Definition

Examples

First Brillouin zone

Properties

### II.2. X-ray and neutron diffraction by a crystal

Diffraction condition

Diffraction by an atom: scattered amplitude

Diffraction by a crystal: structure factor

Symmetry and extinction rules

Beyond basic crystallography: aperiodic crystals and magnetic structures



### II.4. Experiments

How to solve a structure

Technique 1: powder diffraction

Technique 2: single-crystal Laue diffraction

Technique 3: single-crystal (four circle / normal beam) diffraction

## I.4. Beyond basic crystallography: Aperiodic crystals

### Incommensurate modulated structures

Ex: **displacive modulation** along one direction:

$$\vec{r}(\vec{n}, j) = \vec{n} + \vec{r}_j + \vec{U}_j \sin[2\pi \vec{q} \cdot (\vec{n} + \vec{r}_j) + \varphi_j]$$

Average structure

→ 3D lattice

Deviations from the lattice-periodic structure

→ modulation waves

⇒ **(3+d)-dimensional superspace group**

Same approach based on higher-dimension crystallography for:

### Incommensurate composite structures

Subsystems with different incommensurate lattices

### Quasicrystals

No translation symmetry

Presence of non-crystallographic symmetry elements: 5, 8 or 12-fold axes

## I.4. Beyond basic crystallography: Aperiodic crystals

- The six (3+1)-dimensional superspace groups derived from Pnma (ITC, Volume C)

### 9. BASIC STRUCTURAL FEATURES

Table 9.8.3.5. (3 + 1)-Dimensional superspace groups (cont.)

No.	Basic space group	Point group $K_s$	Bravais class No.	Group symbol
62.1	<i>Pnma</i>	( <i>mmm</i> , $1\bar{1}\bar{1}$ )	9	<i>Pnma</i> (00 $\gamma$ )
62.2			9	<i>Pnma</i> (00 $\gamma$ )0s0
62.3			9	<i>Pbnm</i> (00 $\gamma$ )
62.4			9	<i>Pbnm</i> (00 $\gamma$ )s00
62.5			9	<i>Pmcn</i> (00 $\gamma$ )
62.6			9	<i>Pmcn</i> (00 $\gamma$ )s00

- The two icosahedral point groups (quasicrystals):  
$$\begin{cases} 235 & \text{(order 60)} \\ \frac{2}{m} \bar{3}\bar{5} & \text{(order 120)} \end{cases}$$

## II.2. Diffraction: Beyond basic crystallography

### Diffraction by an aperiodic crystal

$$\vec{H} = h\vec{a}^* + k\vec{b}^* + l\vec{c}^* + m_1\vec{q}_1 + m_2\vec{q}_2 \dots + m_d\vec{q}_d$$

Average structure → main reflections:  $\forall i, m_i = 0$

Modulation → satellite reflections of weaker intensity (intensity ↓ when  $m_i \uparrow$ ):  $\forall i, m_i \neq 0$

Reflection conditions for the six (3+1)-dimensional superspace groups derived from *Pnma* (ITC, Volume C)

#### 9. BASIC STRUCTURAL FEATURES

Table 9.8.3.5. (3 + 1)-Dimensional superspace groups (cont.)

No.	Basic space group	Point group $K_s$	Bravais class No.	Group symbol	Special reflection conditions
62.1	<i>Pnma</i>	( <i>mmm</i> , $11\bar{1}$ )	9	<i>Pnma</i> (00 $\gamma$ )	$0klm: k+l=2n; hk00: h=2n$
62.2			9	<i>Pnma</i> (00 $\gamma$ ) $s0s0$	$0klm: k+l=2n; h0lm: m=2n; hk00: h=2n$
62.3			9	<i>Pbnm</i> (00 $\gamma$ )	$0klm: k=2n; h0lm: h+l=2n$
62.4			9	<i>Pbnm</i> (00 $\gamma$ ) $s00$	$0klm: k+m=2n; h0lm: h+l=2n$
62.5			9	<i>Pmcn</i> (00 $\gamma$ )	$h0lm: l=2n; hk00: h+k=2n$
62.6			9	<i>Pmcn</i> (00 $\gamma$ ) $s00$	$0klm: m=2n; h0lm: l=2n; hk00: h+k=2n$

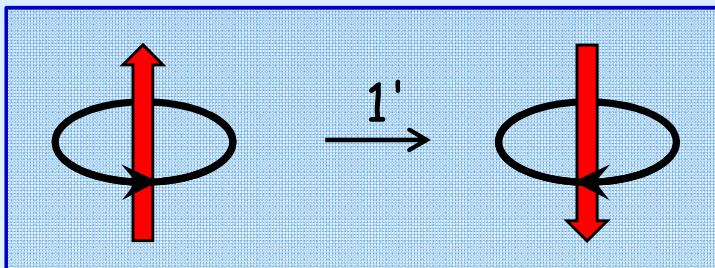
## I.4. Beyond basic crystallography: Magnetic structures

### Invariance properties of magnetic structures

Consider the spin, in addition to the atomic position

Spin = axial vector (magnetic dipole), supposed to be generated by an electrical current

- Additional symmetry operation:  
time reversal = spin reversal =  $1'$
- Rotation: same effect as for polar vector  
(electrical dipole)



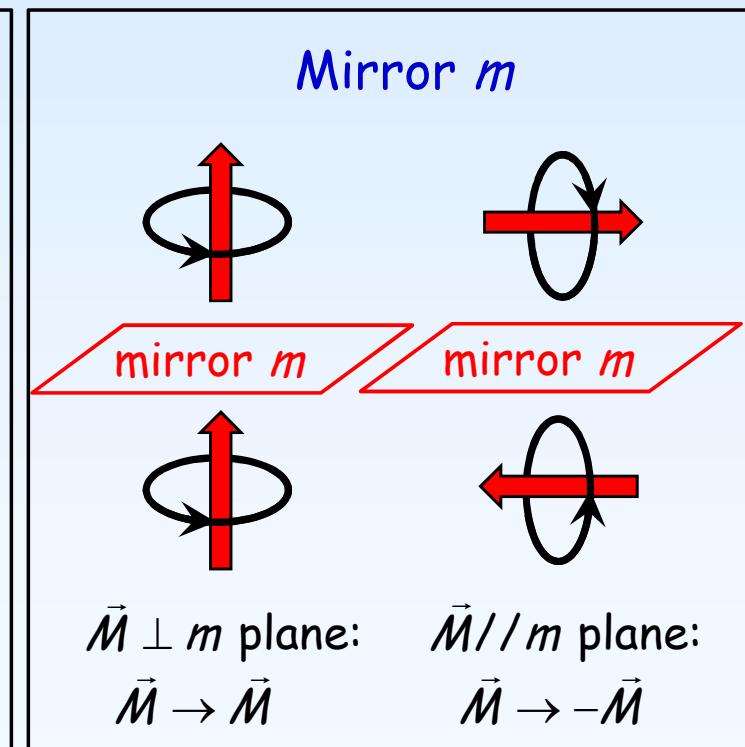
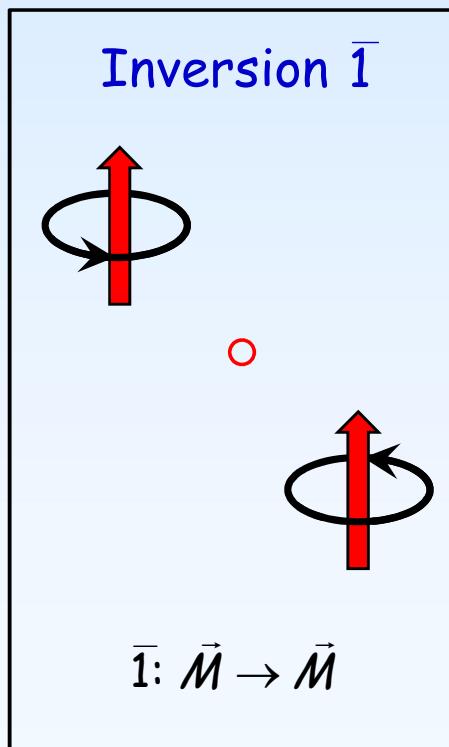
↓

Primed operators  
(= anti-symmetry operators):  
operator  $n$  or  $\bar{n}$  combined with  $1'$

$$\text{e.g.: } m' = m \times 1'$$

$$4' = 4 \times 1'$$

But opposite effect for rotoinversion



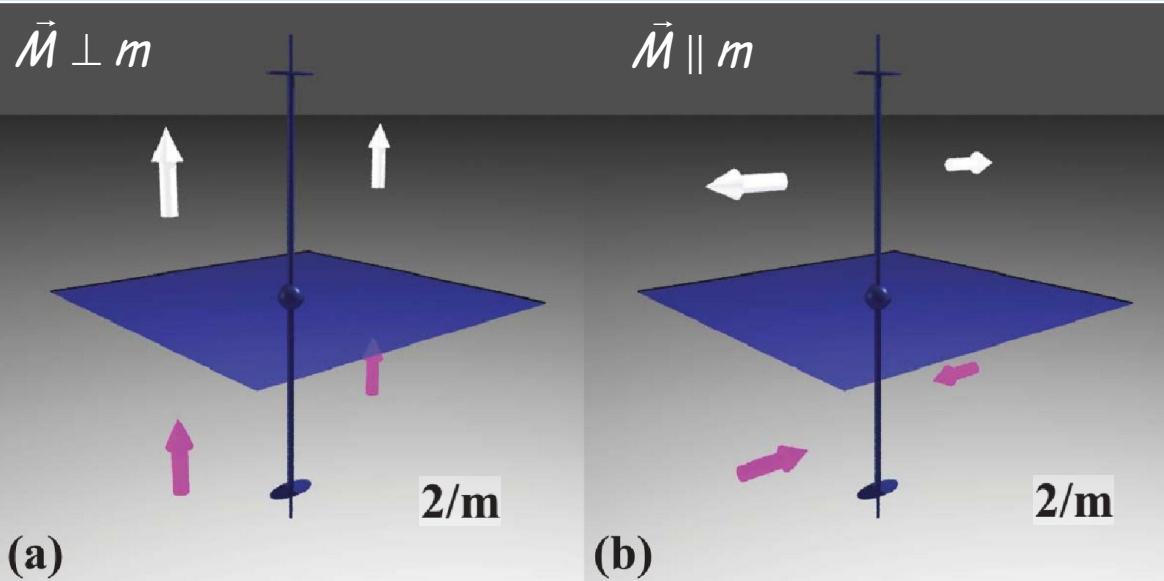
## I.4. Beyond basic crystallography: Magnetic point groups

Ex: Magnetic point groups  $M$  derived from crystallographic point group  $G = 2/m$

1- Colorless magnetic (trivial) PG:

$$M = G$$

$$2/m$$



2- Black-white PG:

$$M = H + (G - H) 1'$$

$2/m =$  admissible magnetic point group  
with admissible spin direction // 2-axis

- conserves a spin located at the origin
- FM state possible

## I.4. Beyond basic crystallography: Magnetic point groups

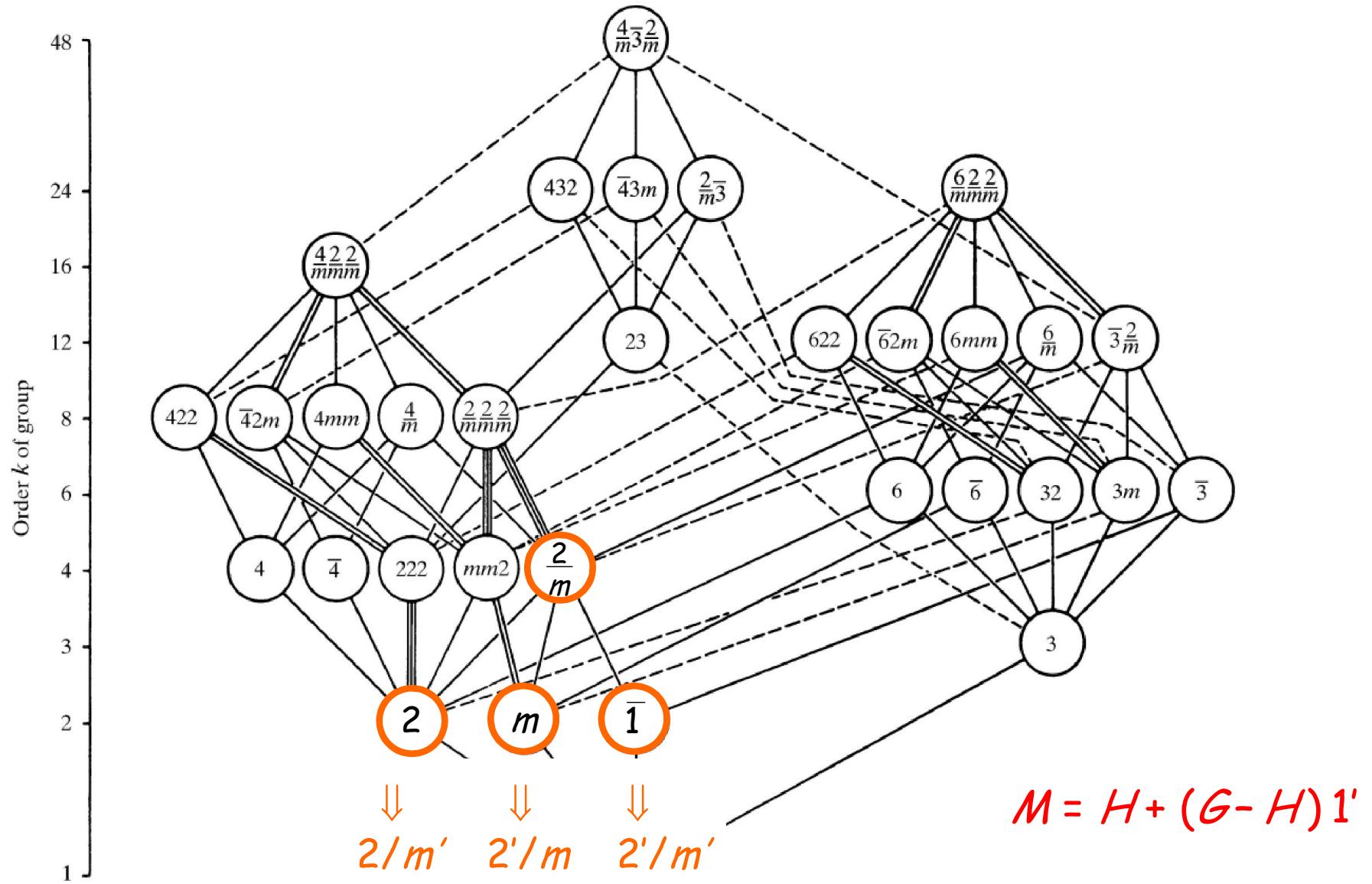


Fig. 10.1.3.2. Maximal subgroups and minimal supergroups of the three-dimensional crystallographic point groups. Solid lines indicate maximal normal subgroups; double or triple solid lines mean that there are two or three maximal normal subgroups with the same symbol. Dashed lines refer to sets of maximal conjugate subgroups. The group orders are given on the left. Full Hermann–Mauguin symbols are used.

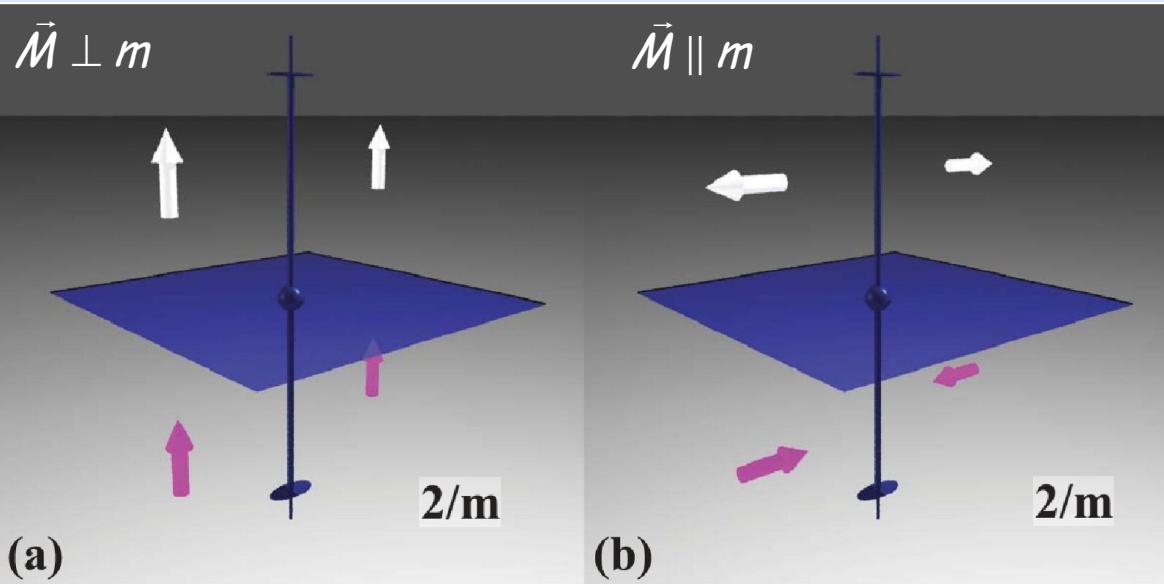
## I.4. Beyond basic crystallography: Magnetic point groups

Ex: Magnetic point groups  $M$  derived from crystallographic point group  $G = 2/m$

1- Colorless magnetic (trivial) PG:

$$M = G$$

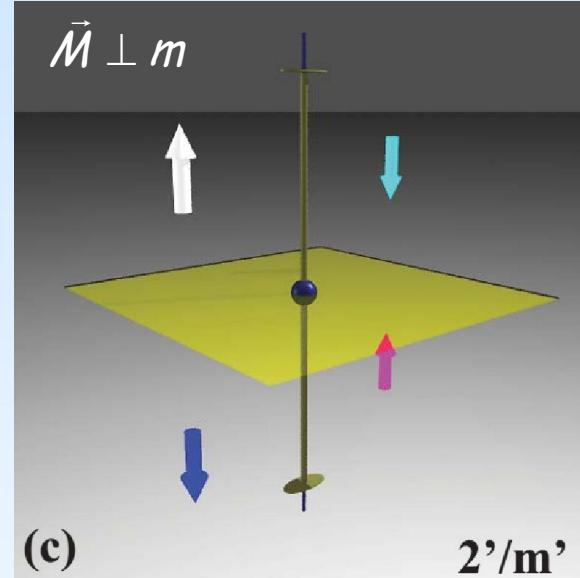
$$2/m$$



2- Black-white PG:

$$M = H + (G - H) 1'$$

$$2/m' \\ 2'/m \\ 2'/m'$$



$2/m =$  admissible magnetic point group  
with admissible spin direction // 2-axis  
→ conserves a spin located at the origin  
→ FM state possible

$2'/m'$  = admissible magnetic PG  
with admissible spin direction  $\perp$  2-axis

3- Gray (paramagnetic) PG:  $2/m1'$

$$M = G + G1'$$

## I.4. Beyond basic crystallography: Magnetic point groups

### The 122 magnetic point groups

32 colorless magnetic point groups (among which 13 admissible)  
58 black-white magnetic point groups (among which 18 admissible)  
32 gray point groups (among which 0 admissible)

### The 31 admissible magnetic point groups

Admissible magnetic point groups			Admissible spin (magnetic moment) direction
1	$\bar{1}$		Any direction
$2'$	$2'/m'$	$m'm2'$	Perpendicular to the 2-fold axis (& to $m$ for $m'm2'$ )
$m'$			Any direction within the plane
$m$			Perpendicular to the plane
$m'm'm$			Perpendicular to the unprimed plane
$2'2'2$			Along the unprimed axis
2	$2/m$	$m'm'2$	Along the 2-fold axis
4	$\bar{4}$	$4/m$	Along the four-fold axis
$4m'm'$	$\bar{4}2m'$	$4/mm'm'$	Along the four-fold axis
3	$\bar{3}$	$32'$	Along the three-fold axis
6	$\bar{6}$	$6/m$	Along the six-fold axis
$6m'm'$	$\bar{6}m'2'$	$6/mm'm'$	Along the six-fold axis

# I.4. Beyond basic crystallography: Magnetic space groups

## The 1651 Shubnikov magnetic space groups

230 colorless trivial magnetic SG of the form  $M = G$

1191 black-white magnetic SG of the form  $M = H + (G - H) 1'$

{ 674 in which the subgroup  $H$  has the same translation as  $G$  (first kind: BW1)  
517 in which the subgroup  $H$  contains anti-translations (second kind: BW2)

230 gray magnetic space groups of the form  $M = G + G1'$

**The magnetic space groups derived from the Fedorov space group: Ima2 (#46)**

Listed with respect to the BNS setting: Belov-Neronova-Smirnova

- #46.241 Ima2 [OG: Ima2 #46.1.338] Type I (Fedorov)
- #46.242 Ima21' [OG: Ima21' #46.2.339] Type II (grey group)
- #46.243 Im'a2' [OG: Im'a2' #46.3.340] Type III (translationgleiche)
- #46.244 Ima'2' [OG: Ima'2' #46.4.341] Type III (translationgleiche)
- #46.245 Im'a2' [OG: Im'a2' #46.5.342] Type III (translationgleiche)
- #46.246 I<sub>c</sub>ma2 [OG: C<sub>1</sub>m'm2' #35.12.247] Type IV (klassengleiche)
- #46.247 I<sub>a</sub>ma2 [OG: A<sub>1</sub>m'm2' #38.13.277] Type IV (klassengleiche)
- #46.248 I<sub>b</sub>ma2 [OG: A<sub>1</sub>bm2' #39.8.285] Type IV (klassengleiche)

Listed with respect to the OG setting: Opechowski-Guccione

- #46.1.338 Ima2 [BNS: Ima2 #46.241] Type I (Fedorov)
- #46.2.339 Ima21' [BNS: Ima21' #46.242] Type II (grey group)
- #46.3.340 Im'a2' [BNS: Im'a2' #46.243] Type III (translationgleiche)
- #46.4.341 Ima'2' [BNS: Ima'2' #46.244] Type III (translationgleiche)
- #46.5.342 Im'a2' [BNS: Im'a2' #46.245] Type III (translationgleiche)
- #46.6.343 I<sub>P</sub>ma2 [BNS: P<sub>1</sub>ma2 #28.98] Type IV (klassengleiche)
- #46.7.344 I<sub>P</sub>m'a2' [BNS: P<sub>1</sub>na2<sub>1</sub> #33.155] Type IV (klassengleiche)
- #46.8.345 I<sub>P</sub>ma2' [BNS: P<sub>1</sub>mc2<sub>1</sub> #26.77] Type IV (klassengleiche)
- #46.9.346 I<sub>P</sub>m'a2' [BNS: P<sub>1</sub>nc2 #30.122] Type IV (klassengleiche)

**bilbao crystallographic server**

[ The crystallographic site at the Condensed Matter Physics Dept. of the University of the Basque Count

[ Space Groups ] [ Layer Groups ] [ Rod Groups ] [ Frieze Groups ] [ Wyckoff Sets ]

<http://www.cryst.ehu.es>

## II.2. Diffraction: Beyond basic crystallography

### Nuclear neutron diffraction

Diffracted intensity

$$I_M(\vec{Q}) \propto |F_N(\vec{Q})|^2 \delta(\vec{Q} - \vec{\tau})$$

with  $F_N(\vec{Q}) = \sum_{j=1}^N b_j e^{2i\pi \vec{Q} \cdot \vec{r}_j} e^{-W_j}$

## II.2. Diffraction: Beyond basic crystallography

### Magnetic neutron diffraction

Diffracted intensity

$$I_M(\vec{Q}) \propto |\bar{F}_{M\perp}(\vec{Q})|^2 \delta(\vec{Q} - \vec{\tau} - \vec{k})$$



Magnetic structure factor  $\bar{F}_M(\vec{Q})$

related to the **magnetic motif**  
(arrangement and values of the magnetic moments in the unit cell)  
⇒ governs the **amplitudes of diffraction**

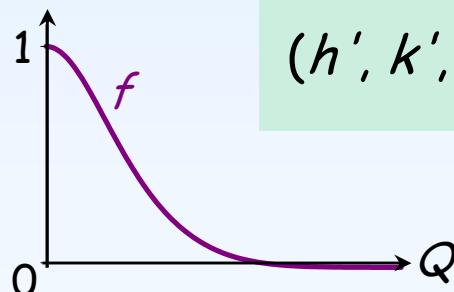
Dirac function  $\delta(\vec{Q} - \vec{\tau} - \vec{k})$

related to the **magnetic lattice**  
(periodicity of the magnetic structure)  
⇒ governs the **directions of diffraction**

$$\bar{F}_M(\vec{Q}) = \sum_{j=1}^{N_{mag}} p f_j(\vec{Q}) \bar{M}_j^k e^{2i\pi \vec{Q} \cdot \vec{r}_j}$$

$\equiv b_j$

$f_j(\vec{Q})$  = Magnetic form factor



$\bar{M}_j^k$  = Fourier component of the magnetic moment

diffraction peaks  $h'k'l'$   
( $h', k', l'$  are not necessarily integers)

## II.2. Diffraction: Beyond basic crystallography

One measures  $|\vec{F}_{M\perp}|^2$  and not  $|\vec{F}_M|^2$

$\vec{F}_{M\perp}$  = projection of  $\vec{F}_M$  in the plane  $\perp \vec{Q}$   
 $(\vec{F}_M \parallel \vec{M})$

$\Rightarrow$  the neutrons are sensitive  
 to the projection of the magnetic moment  
 in the plane  $\perp$  to the scattering vector

if  $\vec{M} \parallel \vec{Q} \Rightarrow F_{M\perp} = 0$

if  $\vec{M} \perp \vec{Q} \Rightarrow F_{M\perp} = F_M$  (i.e. maximum)

Propagation vector  $\vec{k}$

(analogous to the wave vector of a plane wave)

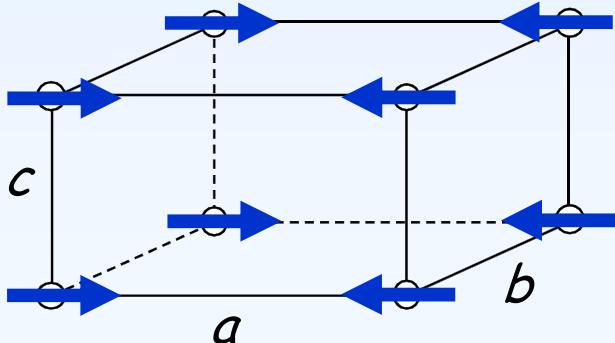
It reflects:

- 1- the periodicity  $L$  of the magnetic structure  
 $(k=0$  if it is the same as the nuclear one,  
 $k=2\pi/L$  otherwise),
- 2- its direction of propagation.

Examples:

- All ferromagnets, some antiferromagnets:  
 $\vec{k} = \vec{0} \rightarrow$  Magnetic peaks at  $\vec{Q} = \vec{\tau} = (h, k, l)$

- AF with periodicity doubled along 1, 2 or 3D:



$$\vec{k} = (\frac{1}{2} \ 0 \ 0)$$

$$\rightarrow \text{Magnetic peaks at } \vec{Q} = \vec{\tau} + \vec{k} = \left( h + \frac{1}{2}, k, l \right)$$

- Sine wave modulated structure along  $c$  (periodicity  $L$ ):

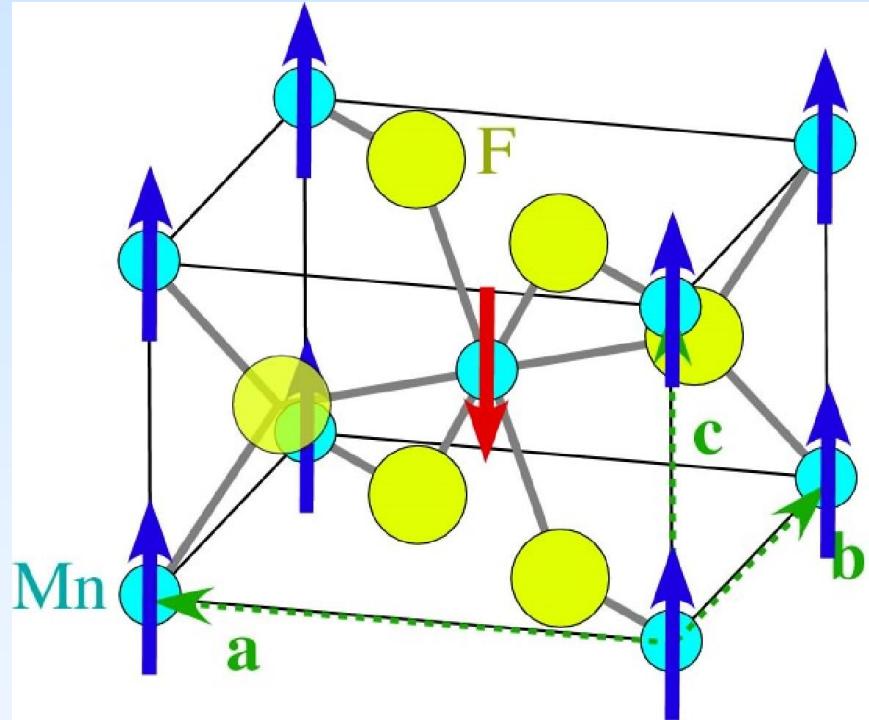
$$\vec{k} = (0 \ 0 \ \frac{1}{L})$$

$$\rightarrow \text{Magnetic peaks at } \vec{Q} = \vec{\tau} \pm \vec{k} = \left( h, k, l \pm \frac{1}{L} \right)$$

## II.2. Diffraction: Beyond basic crystallography

$\text{MnF}_2$ : space group  $P4_2/mnm$

No extinction on  $hkl$

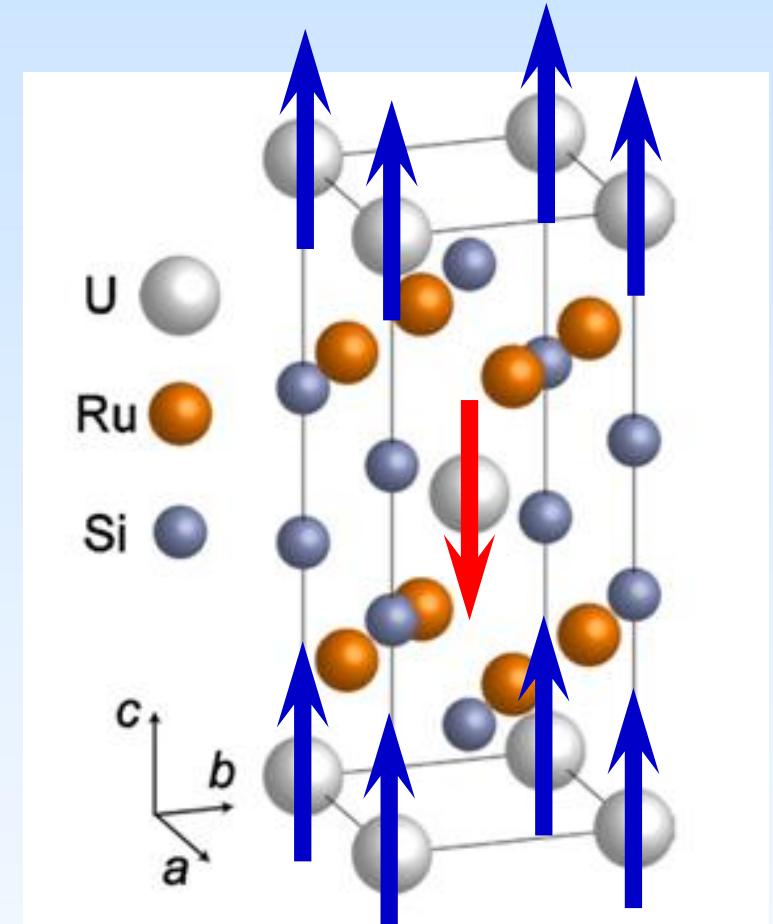


Antiferromagnet with  $\mathbf{k} = (0, 0, 0)$

$$\vec{F}_M(\vec{Q}) = p f_0(\vec{Q}) \vec{M}_0 \left[ 1 - e^{i\pi(h+k+l)} \right] \neq 0 \quad \text{if } h+k+l = 2n+1$$

$\text{URu}_2\text{Si}_2$ : space group  $I4/mmm$

Reflection condition  $hkl$ :  $h+k+l = 2n$



Antiferromagnet with  $\mathbf{k} = (0, 0, 1)$